

Class G Ideal Gas Law + Thermodynamics

G1. Ideal Gas Law

$$\left. \begin{array}{l} \text{Boyle } P \sim \frac{1}{V} \\ \text{Charles } V \sim T \\ \text{Gay-Lussac } P \sim T \end{array} \right\} \frac{P_1 V_1}{T_1} = \frac{P_2 V_2}{T_2}$$

$$P_1 V_1 = P_2 V_2 \quad \Delta T = 0$$

$$\frac{V_1}{T_1} = \frac{V_2}{T_2} \quad \Delta P = 0 \quad T \text{ (Kelvin)}$$

$$\frac{P_1}{T_1} = \frac{P_2}{T_2} \quad \Delta V = 0 \quad \Delta T = 0 \text{ absolute zero}$$

$$PV = nRT$$

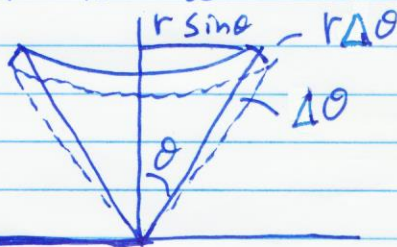
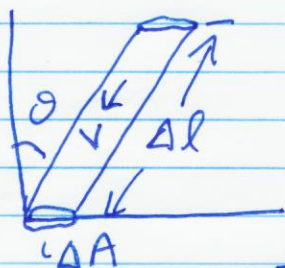
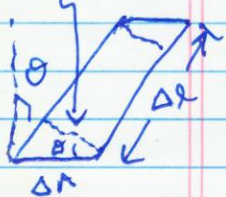
$$PV = NkT$$

moles $n \equiv \frac{N}{N_A}$ ← # particles

N_A ← Avogadro's No. 6.022×10^{23}

$\Delta A \cos \theta$

G2. "Derivation" of the Ideal Gas Law



Bit of pressure

$F = \frac{\Delta P}{\Delta t}$ — momentum $p = mv$
Involves number $P = \frac{F}{A}$

Total $\Delta P = \frac{N}{V} (\Delta A \Delta L \cos \theta) \left[\frac{1}{2} f(v) \Delta v \right] \left[\frac{2\pi r \sin \theta r \Delta \theta}{2\pi r^2} \right] \left[\frac{2mv \cos \theta}{\Delta t} \right] \frac{1}{\Delta A}$

down $\frac{1}{2}$ going up
Prob. v to $v + \Delta v$
Note: $\frac{\Delta L}{\Delta t} = v$

(hemisphere fraction between $\theta + \theta + \Delta \theta$)
Force ← Area
Change in momentum is $2mv \cos \theta$

Note: Quantum Mechanics later teaches us that the lowest energy state has minimal motion.

$$\Delta P = \frac{N}{V} \frac{\Delta A}{\Delta A} \frac{\Delta L}{\Delta t} \frac{1}{2} f(v) \Delta v \cos^2 \theta \sin \theta \Delta \theta 2mv$$

Temperature $T \sim \frac{1}{2} m \overline{v^2}$

average kinetic energy of particles
This definition gives an absolute scale.
No Motion $\Rightarrow T = 0$
"classical"

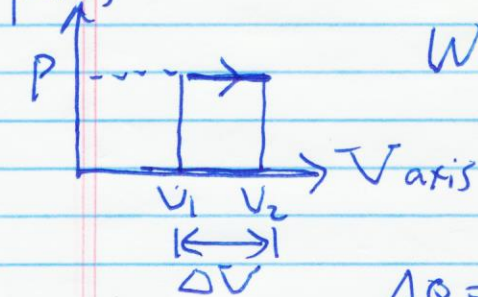
No Relativity (could replace with c)
Speed of light
 $P = \frac{N}{V} m \int_0^\infty v^2 f(v) dv \int_0^{\pi/2} \cos^2 \theta \sin \theta d\theta$

$PV = \frac{N}{3} m \overline{v^2}$ Define T so $PV = NkT$
 $kT = \frac{1}{3} m \overline{v^2}$ $\frac{1}{2} m \overline{v^2} = \frac{3}{2} kT$

$\frac{1}{2} m \overline{v_x^2} = \frac{1}{2} kT$ ← $\frac{1}{2} kT$ for each dimension or degree of freedom
 $\frac{1}{2} m \overline{v^2} = \frac{3}{2} kT$ ← Equipartition Theorem

Change in Energy $\Delta U = \Delta Q - \Delta W$
 Heat Gains in ΔQ
 Work Done by expansion ΔW

G3. First Law of Thermodynamics



Work $\Delta W = F \Delta x = P A \Delta x = P \Delta V$

$U = N \frac{1}{2} m \overline{v^2} = \frac{3}{2} N k T = \frac{3}{2} n R T$
 $n R$

$\Delta Q = \Delta U + P \Delta V$ $\Delta Q_v = \Delta U$
 Specific heat at constant volume constant volume

$\hookrightarrow C_v = \frac{1 \Delta Q_v}{n \Delta T} = \frac{1}{n} \frac{\Delta U}{\Delta T} = \frac{1}{n} \frac{3}{2} n R \frac{\Delta T}{\Delta T} = \frac{3}{2} R$

definition

G4. Four Thermodynamic Processes

1. Isometric $\Delta V = 0$ $W = \int P dV = 0$

2. Isobaric $\Delta P = 0$ $W = \int_{V_1}^{V_2} P dV = P(V_2 - V_1) = P \Delta V$

3. Isothermal $\Delta T = 0$

$W = \int_{V_1}^{V_2} P dV = \int_{V_1}^{V_2} \frac{nRT}{V} dV = nRT \int_{V_1}^{V_2} \frac{1}{V} dV = nRT \ln V \Big|_{V_1}^{V_2}$

$W = nRT \ln \frac{V_2}{V_1}$

$\Delta U = \Delta Q - P \Delta V$ 4. Adiabatic $\Delta Q = 0$ (zero)

$\Delta Q = \Delta U + P \Delta V = n C_v \Delta T + P \Delta V$ (zero) $\Delta Q = n C_v \Delta T + n R \Delta T - V \Delta P$
 $PV = nRT \Rightarrow P \Delta V + V \Delta P = n R \Delta T$ (from) $\Delta Q = n C_p \Delta T$
 $\hookrightarrow n C_v \Delta T = -P \Delta V$ $\hookrightarrow n C_p \Delta T = V \Delta P$

Let $\gamma \equiv \frac{C_p}{C_v}$ $\frac{n C_p \Delta T}{n C_v \Delta T} = \frac{-V \Delta P}{P \Delta V} \Rightarrow \gamma = - \frac{V \Delta P}{P \Delta V}$ $\frac{dP}{P} = -\gamma \frac{dV}{V}$

$\int_{P_1}^{P_2} \frac{dP}{P} = -\gamma \int_{V_1}^{V_2} \frac{dV}{V}$ $\ln P \Big|_{P_1}^{P_2} = -\gamma \ln V \Big|_{V_1}^{V_2}$ $\ln \frac{P_2}{P_1} = -\gamma \ln \frac{V_2}{V_1} = \ln \left[\frac{V_2}{V_1} \right]^{-\gamma}$
 $\frac{P_2}{P_1} = \left(\frac{V_2}{V_1} \right)^{-\gamma}$ $\frac{P_2}{P_1} = \left(\frac{V_1}{V_2} \right)^\gamma$ $P_1 V_1^\gamma = P_2 V_2^\gamma = \text{const}$