

April 7, 2020

U-1

Class U. Green's Functions

U1. Impulse Response

Recall Radioactive Dumping. Let $n(t)$ be number left.

$$\frac{dn(t)}{dt} = -\lambda n(t) + f(t)$$

L dumping (adding more)
L minus due to decay $\sim n(t)$

$$f(t) = \frac{dn(t)}{dt} + \lambda n(t)$$

Take Laplace Transform

$$L\{f(t)\} = L\left\{\frac{dn(t)}{dt}\right\} + \lambda L\{n(t)\}$$

$sN(s) - n(0)$ $N(s)$

$$F(s) = sN(s) + \lambda N(s)$$

$$F(s) = N(s)[s + \lambda]$$

L zero (no waste at the site when we start)

$$N(s) = \frac{F(s)}{s + \lambda} \Rightarrow N(s) = F(s)G(s) \quad G(s) = \frac{1}{s + \lambda}$$

Remember: Product of Laplace Transforms $F(s)G(s) = N(s)$

$\Rightarrow n(t)$ is convolution of $f(t), g(t) \Rightarrow f(t) * g(t)$

$$n(t) = \int_0^t f(u)g(t-u)du$$

All the secrets of the system reside in $G(s), g(t)$

← Isolating $G(s)$ $N(s) = F(s)G(s)$
To find the simplest $n(t)$ L we want $F(s) = 1$

Then $N(s) = 1 \cdot G(s) \Rightarrow N(s) = G(s)$ Simplest!

What $f(t)$ has a Laplace transform of 1?

$$N(s) = F(s)G(s) \quad L\{f(t)\} = F(s) = 1$$

Remember the Sifting Property of the delta function?

$$L \int_0^{\infty} f(t)e^{-st} dt = 1$$

L $\delta(t)$ since $e^{-st}/t=1$ at $t=0$

$$\int \delta(x-a)h(x)dx = h(a)$$

$$N(s) = G(s) = \frac{1}{s + \lambda}$$

$$\int \delta(x)h(x)dx = h(0)$$

From Tables: $h(t) = g(t) = e^{-\lambda t}$

$f(t) = \delta(t)$ is an impulse

$$\text{with } f(t) = A_0 \delta(t) \Rightarrow n(t) = g(t) = A_0 e^{-\lambda t}$$

Pristine empty site — impulse dump at $t=0$ Radioactive Decay

George Green (1793-1841) Self-Educated!!! U-2
 British Mathematical Physicist = Theoretical Physicist

U2. The Green's Function

$D[x(t)] = 0$ Differential Eq. without external input like our dumping function
 $D[x(t)] = S(t)$ Now apply an impulse

Solution will be $g(t)$ Green's function

$$x(t) = \int_0^t f(u) g(t-u) du$$

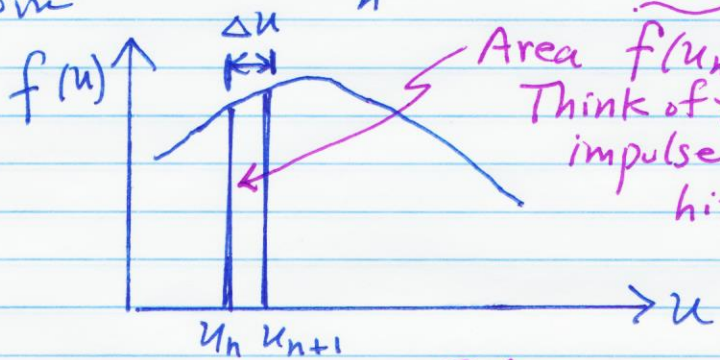
Physics Notation $\int_0^t f(u) G(t,u) du$ \leftarrow Some arbitrary applied function like our dumping or a force in mechanics
 $\int_0^t f(u) G(t,u) du$ \leftarrow $g(t-u)$ by definition for G

More Physics Notation $\int_0^t G(t,u) f(u) du$ \leftarrow Laplace Transform
 $x(t) = \int_0^t G(t,t') f(t') dt'$

\leftarrow for radioactive decay
 $G(t,t') = e^{-(t-t')}$

Remember our discrete sum

$$x(t) = \sum_n G(t, u_n) f(u_n) \Delta u$$



Area $f(u_n) \Delta u$
 Think of this strip as an impulse where with hit the area with the Green's function.

$$G(t, u_n) = e^{-(t-u_n)}$$

The radioactive decay kicks in at time u_n so in the future at some time t you have $e^{-(t-u_n)} f(u_n) \Delta u$ left for that dumping at Δu
 Impulse result applied to the amount dumped in Δu time.

U3. Fourier Transform Space

Recall from our Fourier Transform class $F(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt$

$$F(\omega) \equiv \mathcal{F}\{f(t)\} \quad \mathcal{F}\left\{\frac{df(t)}{dt}\right\} = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{df}{dt} e^{-i\omega t} dt$$

Product Rule leads to integration by parts.

$$\frac{d}{dt} [f(t) e^{-i\omega t}] = \frac{df}{dt} e^{-i\omega t} - i\omega f(t) e^{-i\omega t}$$

← We want this one.

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{df}{dt} e^{-i\omega t} dt = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{d}{dt} [f(t) e^{-i\omega t}] dt + i\omega \int_{-\infty}^{\infty} f(t) e^{-i\omega t} dt$$

$$\mathcal{F}\left\{\frac{df}{dt}\right\} = \frac{1}{\sqrt{2\pi}} f(t) e^{-i\omega t} \Big|_{-\infty}^{\infty} + i\omega F(\omega)$$

↳ Vanishes at $\pm\infty$ No blow-ups
Boundary Conditions

$$\boxed{\mathcal{F}\left\{\frac{df}{dt}\right\} = i\omega F(\omega) \text{ or } i\omega \mathcal{F}\{f(t)\}}$$

Details here.

$$\mathcal{F}\{f''\} = \mathcal{F}\{g'\}$$

where $g = f'$

$$\boxed{\mathcal{F}\left\{\frac{d^2f}{dt^2}\right\} = -\omega^2 F(\omega)}$$

$$\mathcal{F}\{g'\} = i\omega G(\omega)$$

$$\mathcal{F}\{f'\} = i\omega F(\omega) \quad \underbrace{(i\omega)(i\omega) F(\omega)}_{-\omega^2}$$

U4. Finding Green's Functions

by the method extremely useful in theoretical physics

There will be 4 steps.

We will illustrate the method with the radioactive decay differential equation

Don't count this step.

Step 0. What is the differential Equation

$\eta(t)$ amount left at time t at the dumping site

$$\frac{d\eta(t)}{dt} = -\lambda \eta(t) + f(t)$$

external influence
↳ Dumping
↳ radioactive decay

Step 1. Dirac Delta Function Step

$$\frac{d\eta(t)}{dt} = -\lambda \eta(t) + \delta(t)$$

Pick $f(t)$ to be an impulse

Step 2. Fourier Transform
 Left Side → $\frac{dn(t)}{dt} + \lambda n(t) = S(t)$ ← Right Side
 The System Source from outside

$$\mathcal{F}\left\{\frac{dn(t)}{dt}\right\} + \lambda \mathcal{F}\{n(t)\} = \mathcal{F}\{S(t)\}$$

Now, we are in algebraic ω space.

$$i\omega N(\omega) + \lambda N(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} S(t) e^{-i\omega t} dt$$

We need to get back to t space.

$$N(\omega) [i\omega + \lambda] = \frac{1}{\sqrt{2\pi}}$$

$$N(\omega) = \frac{1}{\sqrt{2\pi}} \frac{1}{\lambda + i\omega}$$

The Solution in Fourier Transform ω space.

Step 3. Inverse Fourier Transform

$$n(t) = \mathcal{F}^{-1}\{N(\omega)\}$$

$$n(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} N(\omega) e^{i\omega t} d\omega$$

Note $+i\omega t$
 Note ω integration

Multiply top + bottom by $-i$

$$n(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi}} \frac{1}{\lambda + i\omega} e^{i\omega t} d\omega$$

$$n(t) = \frac{-i}{2\pi} \int_{-\infty}^{\infty} \frac{e^{i\omega t}}{\omega - \lambda i} d\omega$$

Step 4. Complex Integration

Promote $\omega \rightarrow z$

$$\frac{-i}{2\pi} \oint_{CR} \frac{e^{izt}}{z - \lambda i} dz = \frac{-i}{2\pi} \int_{-\infty}^{\infty} \frac{e^{i\omega t}}{\omega - \lambda i} d\omega = \frac{-i}{2\pi} \int_{CR} \frac{e^{izt}}{z - \lambda i} dz$$

This one is easy to do
 We want this to be zero
 We want this to be zero

$$\frac{-i}{2\pi} 2\pi i \text{Res}\left(\frac{e^{izt}}{z - \lambda i}\right) = e^{izt} \Big|_{z=\lambda i} = e^{-\lambda t} = n(t)$$

+1

Note: The text combines Steps 3+4.

V-5

We need to deal with

$$I_{CR} = \frac{-c}{2\pi} \int_{CR} \frac{e^{izt}}{z - \lambda i} dz \quad \text{Let } z = Re^{i\theta}$$

$$dz = iRe^{i\theta} d\theta$$

$$0 \rightarrow \theta \rightarrow \pi$$

$$I_{CR} = \frac{-c}{2\pi} \int_0^\pi \frac{e^{iRe^{i\theta}t}}{Re^{i\theta} - \lambda i} iRe^{i\theta} d\theta$$

$$I_{CR} = \frac{1}{2\pi} \int_0^\pi \frac{e^{iR(\cos\theta + i\sin\theta)t}}{Re^{i\theta} - \lambda i} R e^{i\theta} d\theta$$

Note: Along the imaginary axis $z = iR$
 $e^{izt} \rightarrow e^{-Rt}$ going up e^{+Rt} going down
 the imaginary axis

Important! } Must have the
 Semicircle above + we
 did that. So we are good.

$$\lim_{R \rightarrow \infty} I_{CR} = \frac{1}{2\pi} \lim_{R \rightarrow \infty} \int_0^\pi \frac{e^{iR(\cos\theta + i\sin\theta)t}}{1 - \frac{\lambda i}{Re^{i\theta}}} d\theta$$

$e^{-R\sin\theta} \rightarrow 0$ Killing everything

after dividing top & bottom by $Re^{i\theta}$

Will vanish

What about $\theta = 0$ $R \rightarrow \infty$
 and $e^{-R\sin\theta}$?

The $\theta = 0$ case is on the horizontal axis and that is included in the integration from $-\infty$ to $+\infty$ along the horizontal.

So we can consider I_{CR} starting at $\theta = 0 + \epsilon$

ϵ small (vanishingly small)

So we are okay.

Step 5. The Green's Function

$$\psi(t) = e^{-\lambda t} = G(t, 0)$$

Time Shifted Green's Function

$$G(t, t') = e^{-\lambda(t-t')}$$

General solution for $f(t)$ source function

$$\psi(t) = \int_0^t G(t, t') f(t') dt'$$

↳ Also called a propagator in physics

The source function $f(t)$ effects propagate into the future as the Green function works on it.

Integrate over all times $t' < t$

↳ The present.

Summary of our Journey

